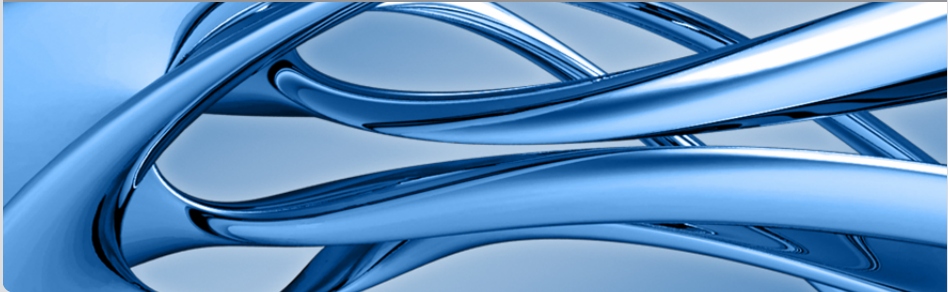


First step towards Parallel and Adaptive Computation of Maxwell's Equations

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For permeability μ and permittivity ε find an electric field \mathbf{E} and a magnetic field \mathbf{H} such that the linear Maxwell system

$$\begin{aligned}\mu\partial_t\mathbf{H} + \nabla\times\mathbf{E} &= \mathbf{0}, & \varepsilon\partial_t\mathbf{E} - \nabla\times\mathbf{H} &= \mathbf{0}, \\ \nabla\cdot(\mu\mathbf{H}) &= 0, & \nabla\cdot(\varepsilon\mathbf{E}) &= 0\end{aligned}$$

holds for all $t \in [0, T]$. For a given initial condition \mathbf{u}_0 this can be written as

$$M\partial_t\mathbf{u}(t) + A\mathbf{u}(t) = \mathbf{0} \quad t \in [0, T], \quad \mathbf{u}(0) = \mathbf{u}_0,$$

where M , A , \mathbf{u} are given by

$$M := \begin{bmatrix} \mu & 0 \\ 0 & \varepsilon \end{bmatrix}, \quad A := \begin{bmatrix} 0 & \nabla\times \\ -\nabla\times & 0 \end{bmatrix}, \quad \mathbf{u} := \begin{bmatrix} \mathbf{H} \\ \mathbf{E} \end{bmatrix}.$$

Consider the 2D reduction

$$\mathbf{u} = (\mathbf{H}_1, \mathbf{H}_2, \mathbf{E}_3), \quad \mathbf{H}_3 \equiv \mathbf{E}_1 \equiv \mathbf{E}_2 \equiv 0$$

for Maxwell's equations in $\Omega \subset \mathbb{R}^2$.

Decomposition for the p -refinement of the space-time domain $Q := \Omega \times (0, T)$:

- Decompose $(0, T)$ such that $[0, T] = \bigcup_{n=0}^{N-1} \bar{I}_n$, where I_n is an open interval $I_n = (t_{n-1}, t_n) \subset (0, T)$ for $n = 1, \dots, N$
- Decompose Ω such that $\bar{\Omega} = \bigcup_K \bar{K}$, where K is an open element (e.g. triangle)
- Let $h = \max \text{diam}(K)$ be the size of K
- Let \mathcal{F}_K be the set of faces of K

Discontinuous Galerkin Method

There exists symmetric matrices $B_1, B_2 \in \mathbb{R}^{3 \times 3}$ so that the linear flux \mathbf{F} is

$$\mathbf{A}\mathbf{u} = \nabla \cdot \mathbf{F}(\mathbf{u}) = B_1 \partial_{x_1} \mathbf{u} + B_2 \partial_{x_2} \mathbf{u}.$$

Multiplying $\mathbf{A}\mathbf{u}$ with a test function \mathbf{v}_K and integrating in K yields

$$\begin{aligned} (\mathbf{A}\mathbf{u}, \mathbf{v}_K)_{0,K} &= \int_K \nabla \cdot \mathbf{F}(\mathbf{u}) \cdot \mathbf{v}_K \, dx \\ &= - \int_K \mathbf{F}(\mathbf{u}) \cdot \nabla \mathbf{v}_K \, dx + \sum_{f \in \mathcal{F}_K} \int_f \mathbf{n}_{K,f} \cdot \mathbf{F}(\mathbf{u}) \cdot \mathbf{v}_K \, da. \end{aligned}$$

Choose p_k , set $V_{K,h} = \mathbb{P}_{p_k}(K)^3$ and define $V_h = \{ \mathbf{v}_h \in L_2(\Omega, \mathbb{R}^3) : \mathbf{v}_h|_K \in V_{K,h} \}$.
 Depending on a numerical flux $\mathbf{n}_K \cdot \mathbf{F}^*(\mathbf{u}_h)$ on $f \in \mathcal{F}_K$ we define $A_h \mathbf{u}_h \in V_h$ by

$$(A_h \mathbf{u}_h, \mathbf{v}_K)_{0,K} = - \int_K \mathbf{F}(\mathbf{u}_h) \cdot \nabla \mathbf{v}_K \, dx + \sum_{f \in \mathcal{F}_K} \int_f \mathbf{n}_{K,f} \cdot \mathbf{F}^*(\mathbf{u}_h) \cdot \mathbf{v}_K \, da$$

for all $\mathbf{u}_h \in V_h$ and $\mathbf{v}_K \in V_{K,h}$ and all K .

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for all $\mathbf{u}_h \in V_h$ and $\mathbf{v}_K \in V_{K,h}$ and all K .

dG Method for Electro-magnetic Waves

For $f \in \mathcal{F}_K$ let K_f be the neighboring tetrahedron with $f = \partial K \cap \partial K_f$.

The outer unit normal on f is denoted by $\mathbf{n}_{K,f}$, and we set $[\mathbf{v}]_{K,f} := \mathbf{v}_{K_f} - \mathbf{v}_K$.

By using the upwind flux, we get the operator

$$\begin{aligned} (A_h(\mathbf{H}_h, \mathbf{E}_h), (\phi_{K,h}, \psi_{K,h}))_{0,K} &= (\nabla \times \mathbf{E}_{K,h}, \phi_{K,h})_{0,K} - (\nabla \times \mathbf{H}_{K,h}, \psi_{K,h})_{0,K} \\ &+ \sum_{f \in \mathcal{F}_K} \left(\frac{c_{K_f \varepsilon_{K_f}}}{c_{K \varepsilon_K} + c_{K_f \varepsilon_{K_f}}} (\mathbf{n}_{K,f} \times [\mathbf{E}_h]_{K,f}, \phi_{K,h})_{0,f} \right. \\ &\quad - \frac{c_{K_f \mu_{K_f}}}{c_{K \mu_K} + c_{K_f \mu_{K_f}}} (\mathbf{n}_{K,f} \times [\mathbf{H}_h]_{K,f}, \psi_{K,h})_{0,f} \\ &\quad + \frac{1}{c_{K \mu_K} + c_{K_f \mu_{K_f}}} (\mathbf{n}_{K,f} \times (\mathbf{n}_{K,f} \times [\mathbf{E}_h]_{K,f}), \psi_{K,h})_{0,f} \\ &\quad \left. + \frac{1}{c_{K \varepsilon_K} + c_{K_f \varepsilon_{K_f}}} (\mathbf{n}_{K,f} \times (\mathbf{n}_{K,f} \times [\mathbf{H}_h]_{K,f}), \phi_{K,h})_{0,f} \right) \end{aligned}$$

for $(\mathbf{E}_h, \mathbf{H}_h) \in V_h$ and $(\psi_{K,h}, \phi_{K,h}) \in V_{K,h}$.

(for details see Hesthaven and Warburton 2002)

Time Integration

Consider the semi-discrete problem

$$M_h \partial_t \mathbf{u}_h(t) + A_h \mathbf{u}_h(t) = \mathbf{0} \quad \text{for } t \in [0, T] \quad \text{subject to } \mathbf{u}_h(0) = \mathbf{u}_{h,0}.$$

To solve this problem, we use the implicit midpoint rule

$$\mathbf{u}_h^n = \mathbf{u}_h^{n-1} + (t_n - t_{n-1}) \left(M_h + \frac{t_n - t_{n-1}}{2} A_h \right)^{-1} A_h \mathbf{u}_h^{n-1} \quad \text{for } n = 1, \dots, N$$

with step size $t_n - t_{n-1}$ and initial condition $\mathbf{u}_h^0 = \mathbf{u}_{h,0}$.

Properties of the implicit midpoint rule:

- No CFL (Courant-Friedrichs-Lewy) condition \Rightarrow allows larger time steps $t_n - t_{n-1}$
- Convergence order 2 in time
- The upwind flux guarantees that $(M_h + \frac{t_n - t_{n-1}}{2} A_h)$ is positive definite
- Costs: solve $(M_h + \frac{t_n - t_{n-1}}{2} A_h) \tilde{\mathbf{u}}_h^n = A_h \mathbf{u}_h^{n-1}$ in each step

Hence the solution \mathbf{u}_h is computed sequentially on the slices $S_n := \Omega \times I_n \subset Q$.

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Define space-time cells

$$\tau = K_\tau \times I_\tau$$

which consist of a spatial element K_τ and a local time interval $I_\tau = (t_\tau^{\min}, t_\tau^{\max})$ and decompose Q into a finite number of open space-time elements $\tau \subset Q$ such that

$$\bar{Q} = \bigcup_{\tau \in \mathcal{T}} \bar{\tau}.$$

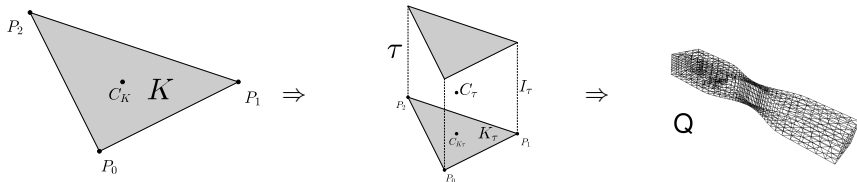


Figure: spatial element, space-time element, space-time mesh

Space-time dG approximation

For every τ choose polynomial degrees p_τ for the ansatz in space, and define the local test spaces $H_{\tau,h} = \mathbb{P}_{p_\tau}(K_\tau)^3$ and the test space

$$H_h = \{ \mathbf{v}_h \in L_2(Q)^3 : \mathbf{v}_h|_\tau \in H_{\tau,h} \}.$$

For the ansatz space, we define the affine space depending on the initial condition

$$U_h = \{ \mathbf{u}_h \in H^1(0, T; L_2(\Omega)^3) : \mathbf{u}_h(0) = \mathbf{u}_0 \text{ and for all } \tau \in \mathcal{T} \text{ and } (\mathbf{x}, t) \in \tau$$

$$\mathbf{u}_h(\mathbf{x}, t) = \frac{t_\tau^{\max} - t}{t_\tau^{\max} - t_\tau^{\min}} \mathbf{w}_{\tau,h}(\mathbf{x}, t_\tau^{\min}) + \frac{t - t_\tau^{\min}}{t_\tau^{\max} - t_\tau^{\min}} \mathbf{v}_{\tau,h}(\mathbf{x}),$$

$$\text{where } \mathbf{w}_{\tau,h} \in U_h|_{[0, t_\tau^{\min}]} \text{ and } \mathbf{v}_{\tau,h} \in H_{\tau,h} \}$$

Let A_h be the discontinuous Galerkin operator with upwind flux approximating A .

Lemma

Let $L_h = M_h \partial_t + A_h$ and $\mathbf{f} \in L_2(Q)^3$. The discrete solution $\mathbf{u}_h \in U_h$ of the implicit midpoint rule is characterized by the variational equation

$$(L_h \mathbf{u}_h, \mathbf{v}_h)_Q = (\mathbf{f}, \mathbf{v}_h)_Q, \quad \mathbf{v}_h \in H_h.$$

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Data Structure

Our implementation is done in C++ and the data structure is organized as follows:

- (Time-)cells, faces, edges are identified by their geometric midpoints
- Hash maps containers are used for the data
- Geometric midpoints are used as hash keys

```
class Cell : public vector<Point> {...};
class Cells : public hash_map<Point,Cell*,Hash> {...};

class Interval {
    const double t_min;
    const double t_max;
    ...};

class TCell {
    const Cell* C;
    const Interval* I;
    ...};

class TCells : public hash_map<Point,TCell*,Hash> {
    tcell tcells () const { return tcell(begin());}
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The use of hash maps allows us to distribute cells among the different processes and solve the problem in parallel. Every cell is stored on only one master process and communicates its data to other processes, if needed.

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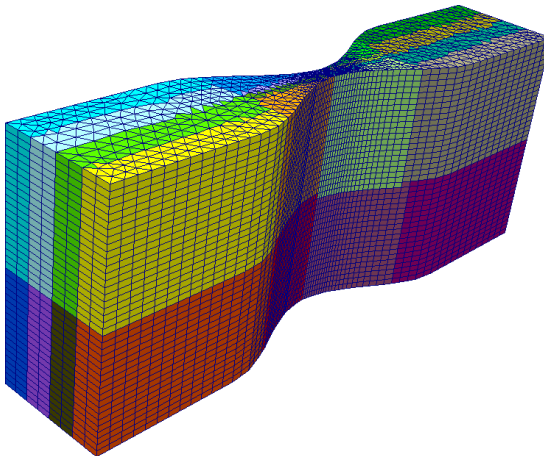


Figure: Distribution of the space-time cells on 32 processes

ρ -adaptive refinement in space:

- Compute \mathbf{u}_h with lowest polynomial degree $p = 0$ on all slices S_n
- Refinement indicator η_τ of a cell τ :

$$\eta_\tau^2 := \sum_{f \in \mathcal{F}_\tau} \eta_f^2 \quad \text{where } \eta_f^2 := h_f \|(\mathbf{F}^*(\mathbf{u}_h) - \mathbf{F}(\mathbf{u}_h)) \cdot \mathbf{n}_f\|_{L_2(f)}^2,$$

with flux and numerical flux \mathbf{F} and \mathbf{F}^* , and area h_f , and the outer normal vector \mathbf{n}_f of the face f .

- Increase the polynomial degree p on the space-time cell $\tau \subset S_n$, if

$$\eta_\tau > (1 - \theta) \max_{\tau \in S_n} \eta_\tau$$

holds for η_τ and a given parameter $\theta \in [0, 1]$, e.g. $\theta = 0.99$.

Space-time refinement:

- Same heuristic as for p -refinement
- If $\tau = K_\tau \times l_\tau$ should be refined in time, then
 1. Split $l_\tau = (t_\tau^{\min}, t_\tau^{\max})$ into l_τ^1 and l_τ^2 , s.t. $\overline{l_\tau} = \overline{l_\tau^1} \cup \overline{l_\tau^2}$,
where $l_\tau^1 := (t_\tau^{\min}, t_\tau^{1/2})$, $l_\tau^2 := (t_\tau^{1/2}, t_\tau^{\max})$ and $t_\tau^{1/2} := 0.5(t_\tau^{\min} + t_\tau^{\max})$
 2. Replace τ by two new space-time cells $\tau^1 = K_\tau \times l_\tau^1$ and $\tau^2 = K_\tau \times l_\tau^2$
- Perform the implicit midpoint rule twice on time refined space-time cells

Performance of the heuristic approach:

- **High** polynomial degrees and **small** time steps are used in areas where a single wavefront is located
- **Lowest** polynomial degrees and **larger** time steps are used in areas with absence of a wave

Space-time refinement:

- Same heuristic as for p -refinement
- If $\tau = K_\tau \times I_\tau$ should be refined in time, then
 1. Split $I_\tau = (t_\tau^{\min}, t_\tau^{\max})$ into I_τ^1 and I_τ^2 , s.t. $\overline{I_\tau} = \overline{I_\tau^1} \cup \overline{I_\tau^2}$,
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We consider:

- Unstructured triangular mesh in a locally tapered domain $\Omega \subset (0, 10) \times (-1, 1)$ with reflecting boundary conditions

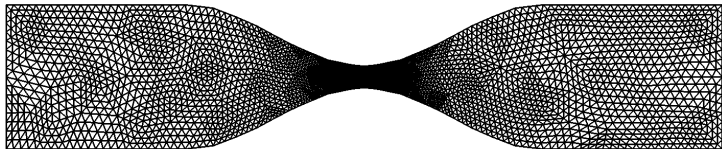


Figure: Unstructured triangular mesh in tapered domain Ω

- Constant parameters $\mu = \varepsilon = 1$
- Initial condition $\mathbf{u}_0(x) = \begin{cases} (0, 0, \cos(4\pi x_1 - 3\pi) + 1)^T, & \text{for } 1 \leq x_1 \leq 1.5, \\ \mathbf{0}, & \text{else.} \end{cases}$
- Final time $T = 8$ with constant initial step size $t_n - t_{n-1} = 0.1$

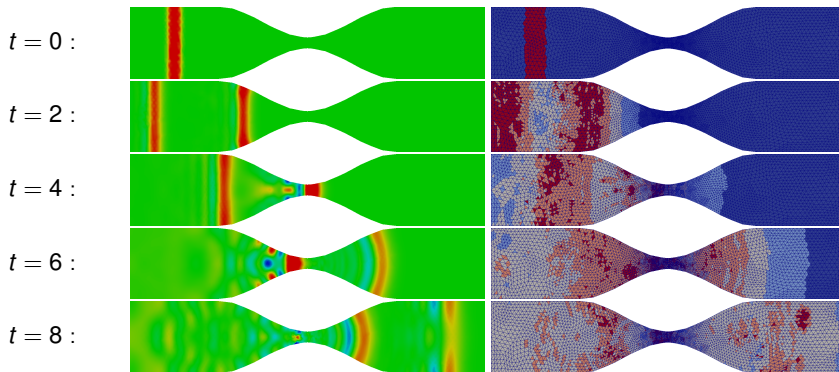


Figure: Initial distribution of \mathbf{E}_3 and polynomial degrees and time evolution

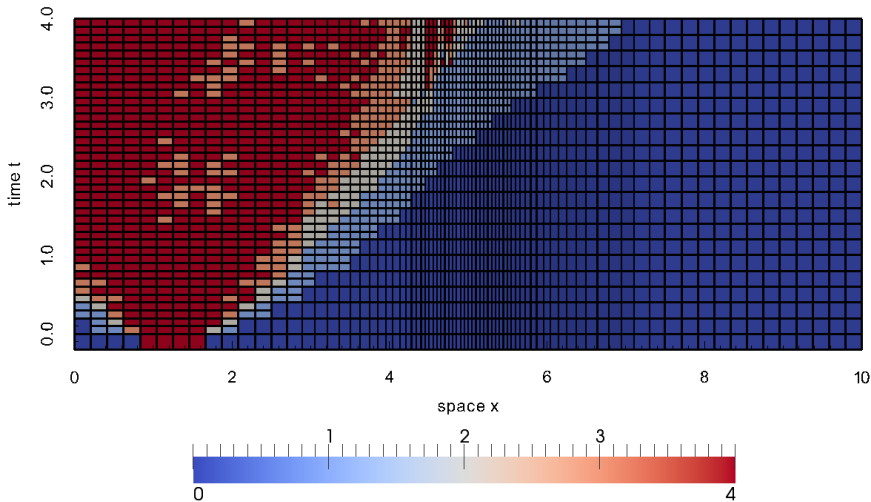


Figure: Polynomial degrees and time refinement in the time interval $[0, 4]$



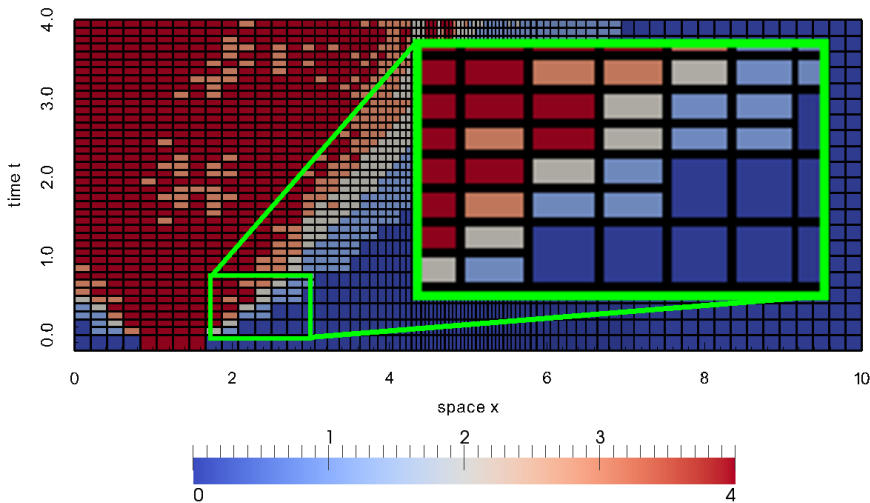


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Adaptive polynomial degrees:

ρ	ρ -adaptive		uniform	
	degrees of freedom	$\ \mathbf{u}_{\rho+1}(T) - \mathbf{u}_{\rho}(T)\ _V$	degrees of freedom	$\ \mathbf{u}_{\rho+1}(T) - \mathbf{u}_{\rho}(T)\ _V$
0	1,469,664		1,469,664	
1	3,303,156	0.835850	4,408,992	0.853796
2	5,424,354	0.269539	8,817,984	0.265763
3	8,100,078	0.086987	14,656,640	0.020959
4	11,494,896	0.022396	22,044,960	0.002958

Table: Adaptive polynomial degrees and corresponding degrees of freedom

Space-time adaptive (max. one refinement in time):

p	p -time-adaptive		uniform	
	degrees of freedom	$\ \mathbf{u}_{p+1}(T) - \mathbf{u}_p(T)\ _V$	degrees of freedom	$\ \mathbf{u}_{p+1}(T) - \mathbf{u}_p(T)\ _V$
0	1,469,664		2,939,328	
1	6,051,378	0.840901	8,817,984	0.845272
2	10,040,952	0.309939	17,635,968	0.289086
3	13,678,458	0.083207	29,393,280	0.023039
4	16,445,328	0.057660	44,089,920	0.003572

Table: Space-time adaptive refinement and corresponding degrees of freedom

- Computation of the full 3D Maxwell problem
- Additional refinements in time
- h -adaptivity in space
- Better error indicators and estimators